



Institut für Mathematik
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computational mathematics

Composite Finite Element Methods Part 1: Finite Element Methods

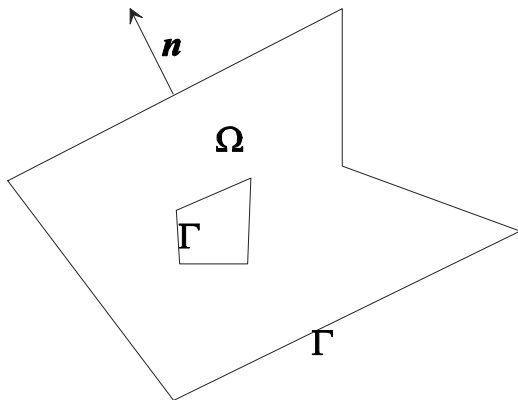
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June 2023

Model Problem:

Let $\Omega \subset \mathbb{R}^d$ be a bounded open interval ($d = 1$), a bounded polygonal domain ($d = 2$), a bounded polyhedral domain ($d = 3$), etc., with boundary $\Gamma := \partial\Omega$.



Polygonal domain Ω
with boundary $\Gamma := \partial\Omega$
and outward normal \mathbf{n} .

Notation

$$\nabla v := \left(\partial_j v \right)_{j=1}^d, \quad \Delta v := \sum_{j=1}^d \partial_j^2 v, \quad \operatorname{div} \mathbf{w} := \sum_{j=1}^d \partial_j w_j.$$

Poisson problem (classical formulation):

Let a “load function” $f : \Omega \rightarrow \mathbb{R}$ be given. Find a function $u : \Omega \rightarrow \mathbb{R}$ such that

$$\begin{aligned} -\Delta u(\mathbf{x}) &= f(\mathbf{x}) & \forall \mathbf{x} \in \Omega, \\ u(\mathbf{x}) &= 0 & \forall \mathbf{x} \in \Gamma. \end{aligned}$$

Remark. This formulation requires the equality in every point $\mathbf{x} \in \Omega$. For existence, and uniqueness results this formulation has severe drawback.

Variational formulation of the Poisson model problem (I)

For the *weak* or *variational formulation* one replaces the pointwise conditions by *integral conditions*. Let $v \in C^\infty(\bar{\Omega})$ be a *test function*. Multiplying the differential equation by v and integrating over the domain Ω leads to:
find $u : \Omega \rightarrow \mathbb{R}$ with $u|_\Gamma = 0$ such that

$$\int_{\Omega} (-\Delta u) v = \int_{\Omega} f v \quad \forall v \in C^\infty(\bar{\Omega}).$$

We integrate by parts to obtain

$$\int_{\Omega} \langle \nabla u, \nabla v \rangle - \int_{\Gamma} \frac{\partial u}{\partial \mathbf{n}} v = \int_{\Omega} f v \quad \forall v \in C^\infty(\bar{\Omega}).$$

Variational formulation of the Poisson model problem (II)

Since the function u is zero on Γ we may restrict to test functions v which are zero on the boundary. Hence, the boundary integral can be dropped and we have derived the weak formulation:

Find $u \in V_0$ such that

$$\int_{\Omega} \langle \nabla u, \nabla v \rangle = \int_{\Omega} f v \quad \forall v \in C_0^{\infty}(\Omega)$$

with $C_0^{\infty}(\Omega) := \{v \in C^{\infty}(\Omega) : v|_{\Gamma} = 0\}$.

Variational formulation of the Poisson model problem (III)

The *energy space* V_0 must satisfy:

- a) V_0 is a Hilbert space,
- b) $a(u, v) := \int_{\Omega} \langle \nabla u, \nabla v \rangle$ defines a scalar product in V_0 ,
- c) functions in V_0 are zero on Γ .

The subspace of $L^2(\Omega)$ with derivatives in $L^2(\Omega)$ is the *Sobolev space*

$$H^1(\Omega) := \left\{ u \in L^2(\Omega) \mid \int_{\Omega} \|\nabla u\|^2 < \infty \right\}.$$

Trace theorem: Functions in $H^1(\Omega)$ have well-defined restrictions to the boundary $\partial\Omega$.

Definition (energy space) The *energy space* for the Poisson model problem is

$$V_0 := \{v \in H^1(\Omega) \mid v|_{\Gamma} = 0\}$$

The **weak (variational) formulation** of the Poisson problem reads:
find $u \in V_0$ such that

$$a(u, v) = F(v) \quad \forall v \in V_0.$$

Here, the bilinear form $a : V_0 \times V_0 \rightarrow \mathbb{R}$ and the functional $F : V_0 \rightarrow \mathbb{R}$ is given by

$$a(u, v) := \int_{\Omega} \langle \nabla u, \nabla v \rangle \quad \text{and} \quad F(v) := \int_{\Omega} f v \quad \forall v \in V_0.$$

Analysis on the continuous problem

Definition. Let V be a Hilbert space with norm $\|\cdot\|_V$ and let a bilinearform $a : V \times V \rightarrow \mathbb{R}$ be given.

$a(\cdot, \cdot)$ is *continuous* if there is some $C > 0$ such that

$$|a(v, w)| \leq C \|v\|_V \|w\|_V \quad \forall v, w \in V.$$

$a(\cdot, \cdot)$ is *coercive* if there is some $c > 0$ such that

$$|a(v, v)| \geq c \|v\|_V^2 \quad \forall v \in V.$$

Definition (cont'd)

$a(\cdot, \cdot)$ is symmetric if

$$a(v, w) = a(w, v) \quad \forall v, w \in V.$$

A linear form $F \in V'$ is *continuous* if

$$\|F\|_{V'} := \sup_{v \in V \setminus \{0\}} \frac{|F(v)|}{\|v\|_V} < \infty.$$

Theorem (Lax-Milgram). Let V be a Hilbert space and $a : V \times V \rightarrow \mathbb{R}$ be symmetric, continuous, and coercive. Then, the variational problem: for given continuous linear form $F \in V'$, find $u \in V$ such that

$$a(u, v) = F(v) \quad \forall v \in V$$

has a unique solution which satisfies

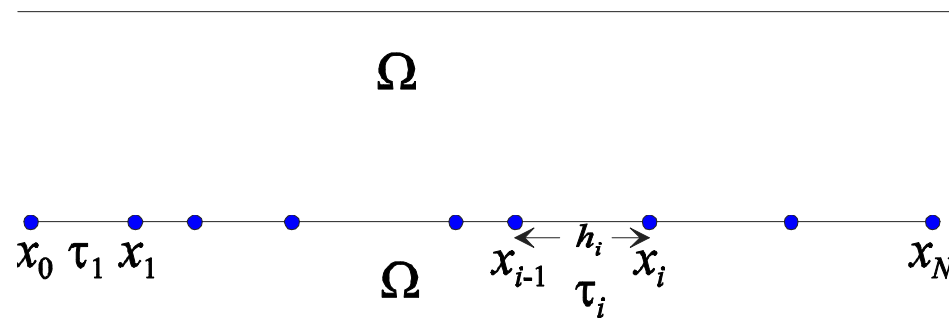
$$\|u\|_V \leq \frac{1}{c} \|F\|_{V'}.$$

Galerkin Finite Element Method:

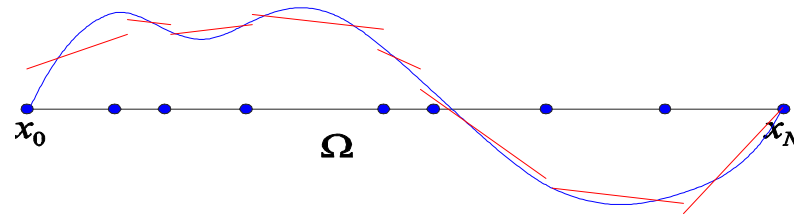
To approximate the continuous problem, a *finite-dimensional* function space $S \subset V_0$ has to be defined.

Idea of finite elements:

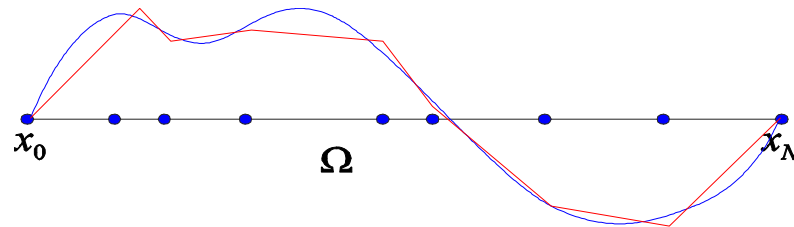
a) subdivide Ω in small *simplices* (intervals, triangles, tetrahedrons)



b) approximate on the simplices by piecewise *polynomials*



c) Enforce continuity across element boundaries and boundary conditions to ensure $S \subset V_0$

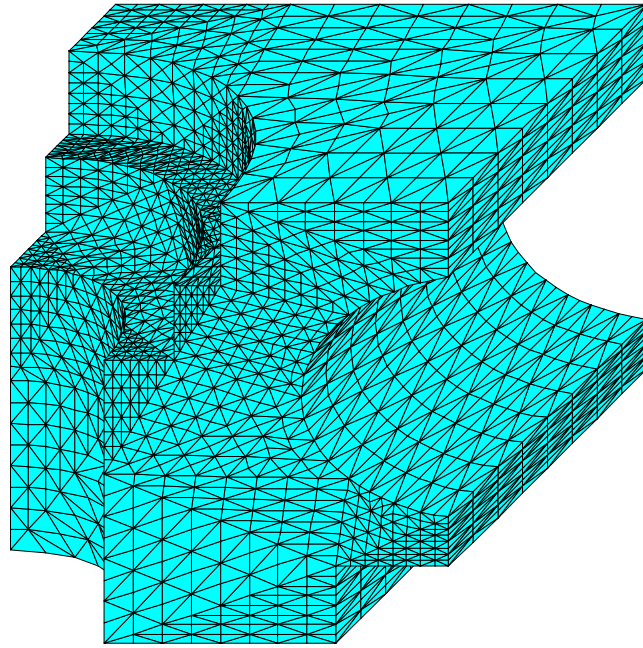


Definition (shape regularity, mesh width). Let $\mathcal{T} := \{K_j : 1 \leq j \leq N\}$ denote a conforming (no hanging nodes), simplicial finite element mesh for Ω . The (local) mesh width is given by

$$h_K := \text{diam } K \quad \text{and} \quad h := \max \{h_K : K \in \mathcal{T}\}.$$

For the approximation quality, the *shape regularity constant* is important

$$\gamma_{\text{sr}}(\mathcal{T}) := \max \left\{ \frac{h_K^d}{|K|} : K \in \mathcal{T} \right\}.$$



Example of a tetrahedral mesh with good shape regularity constant.

Definition (Finite Element space). Let $\mathcal{T} := \{K_j : 1 \leq j \leq N\}$ denote a conforming, simplicial finite element mesh for Ω and $p \geq 1$. Then

$$S_{\mathcal{T}}^p := \{u \in C^0(\Omega) \mid \forall K \in \mathcal{T} : u|_K \in \mathbb{P}_p(K)\}$$
$$S := S_{\mathcal{T},0}^p := \{u \in S_{\mathcal{T}}^p \mid u|_{\partial\Omega} = 0\}.$$

Definition (Galerkin method). The *Galerkin discretization* of a variational problem is characterized by a finite-dimensional subspace $S \subset V_0$, $N := \dim S < \infty$:

Find $u_S \in S$ such that

$$a(u_S, v) = F(v) \quad \forall v \in S.$$

Stability and convergence analysis

Theorem (Céa). Let V be a Hilbert space and $a : V \times V \rightarrow \mathbb{R}$ be symmetric, continuous, and coercive. Let $S \subset V$ with $\dim S < \infty$. Then, the Galerkin method has a *unique solution* which satisfies the quasi-optimal error estimate

$$\|u - u_S\|_V \leq \frac{C}{c} \inf_{v \in S} \|u - v\|_V.$$

The *Galerkin orthogonality* holds

$$a(u - u_S, v) = 0 \quad \forall v \in S.$$

Theorem. If the exact solution of the Poisson model problem is *regular*, i.e., $u \in H_0^1(\Omega) \cap H^{p+1}(\Omega)$ then the *energy error* satisfies

$$\|u - u_S\|_{H^1(\Omega)} \leq \frac{C}{c} C_{\text{sr}} h_T^p \|u\|_{H^{p+1}(\Omega)}.$$

Computational aspects:

For the numerical solution, a basis for S is needed

$$S = \text{span} \{B_i : i \in \mathcal{I}\} \quad \text{with} \quad |\mathcal{I}| = \dim S = N.$$

Basis representation of Galerkin discretization

The *stiffness (system) matrix* $\mathbf{A} = (a_{i,j})_{i,j=1}^N \in \mathbb{R}^{N \times N}$ and the *load vector (right-hand side)* $\mathbf{r} := (r_i)_{i=1}^N$ are given by

$$a_{i,j} := a(B_j, B_i) = \int_{\Omega} \langle \nabla B_j, \nabla B_i \rangle,$$
$$r_i = F(B_i) = \int_{\Omega} f B_i.$$

The Galerkin solution u_S has a unique basis representation

$$u_S = \sum_{i \in \mathcal{I}} u_i B_i$$

and the coefficient vector $\mathbf{u} = (u_i)_{i \in \mathcal{I}}$ is the unique solution of the system of linear equations

$$\mathbf{A}\mathbf{u} = \mathbf{r}$$

Remark. The coercivity and symmetry of $a(\cdot, \cdot)$ implies that the matrix \mathbf{A} is *symmetric, positive definite (spd)*.

Affine equivalence

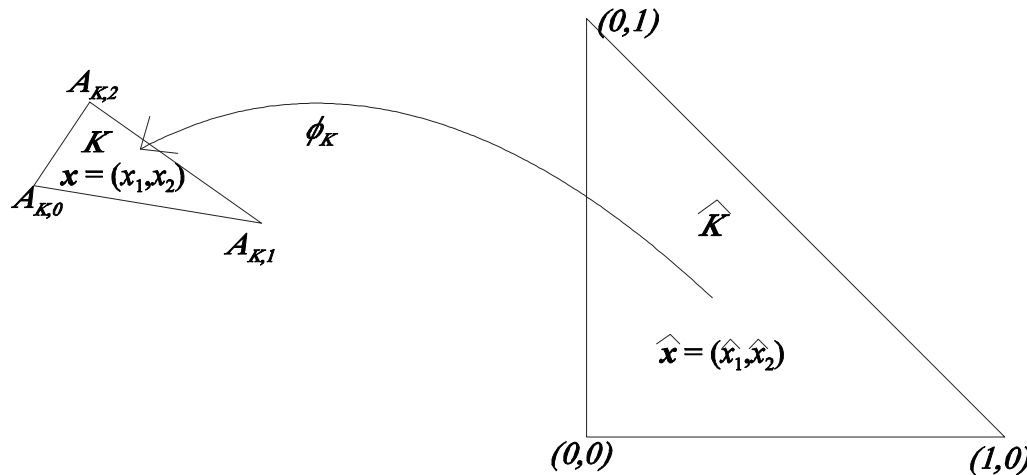
All finite element computations should be transformed to the *affine equivalent* reference element:

$$\widehat{K} := \left\{ \widehat{\mathbf{x}} = (\widehat{x}_i)_{i=1}^d \in \mathbb{R}_{\geq 0}^d \mid \sum_{i=1}^d \widehat{x}_i \leq 1 \right\}.$$

Then any simplex K with vertices $\mathbf{A}_{K,j}$, $0 \leq j \leq d$, has an affine pullback $\phi_K : \widehat{K} \rightarrow K$ given by

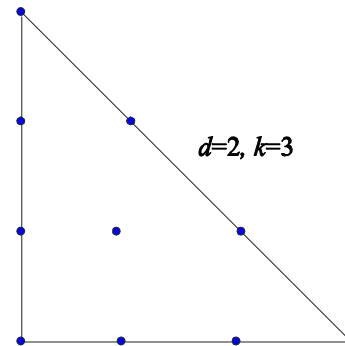
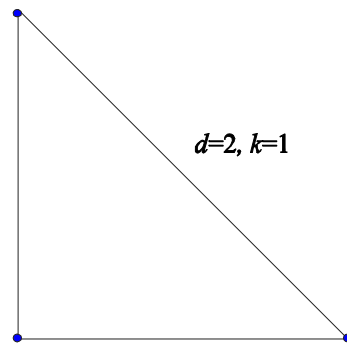
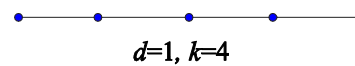
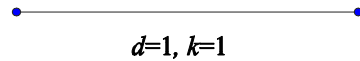
$$\phi_K(\widehat{\mathbf{x}}) = \mathbf{A}_{K,0} + \mathbf{m}_K \widehat{\mathbf{x}}$$

with the $d \times d$ matrix \mathbf{m}_K having column vectors $\mathbf{A}_{K,j} - \mathbf{A}_{K,0}$.



Basis functions of $\mathbb{P}_p(\widehat{K})$ are defined by using nodal points $\widehat{\mathcal{N}}_k$:

$$\widehat{\mathcal{N}}_k := \left\{ \frac{i}{p} : 0 \leq i \leq p \right\}^d \cap \widehat{K}$$



For all $\mathbf{z} \in \mathcal{N}_k$ the *Lagrange* basis function $\hat{B}_{\mathbf{z}} \in \mathbb{P}_k(\widehat{K})$ is characterized by

$$\hat{B}_{\hat{\mathbf{z}}}(\hat{\mathbf{y}}) := \begin{cases} 1 & \hat{\mathbf{y}} = \hat{\mathbf{z}}, \\ 0 & \hat{\mathbf{y}} \in \widehat{\mathcal{N}}_k \setminus \{\hat{\mathbf{z}}\}. \end{cases}$$

$$\left. \begin{aligned} \hat{B}_{(0)}(\hat{x}) &= 1 - \hat{x} \\ \hat{B}_{(1)}(\hat{x}) &= \hat{x} \end{aligned} \right\} d = 1, k = 1$$

$$\left. \begin{aligned} \hat{B}_{(0,0)}(\hat{\mathbf{x}}) &= 1 - \hat{x}_1 - \hat{x}_2 \\ \hat{B}_{(1,0)}(\hat{\mathbf{x}}) &= \hat{x}_1 \\ \hat{B}_{(0,1)}(\hat{\mathbf{x}}) &= \hat{x}_2 \end{aligned} \right\} d = 2, k = 1$$

Global nodal points and basis functions are defined via affine pullbacks:

$$\begin{aligned}\mathcal{N}_k(\mathcal{T}) &:= \left\{ \phi_K(\hat{\mathbf{z}}) : K \in \mathcal{T}, \hat{\mathbf{z}} \in \widehat{\mathcal{N}}_k \right\}, \\ \mathcal{N}_{k,0}(\mathcal{T}) &:= \left\{ \mathbf{z} \in \mathcal{N}_k(\mathcal{T}) \mid \mathbf{z} \notin \partial\Omega \right\}.\end{aligned}$$

For $\mathbf{z} \in \mathcal{N}_k(\mathcal{T})$ the *nodal patch* is given by

$$\begin{aligned}\mathcal{T}_{\mathbf{z}} &:= \{K \in \mathcal{T} \mid \mathbf{z} \in K\} \\ \omega_{\mathbf{z}} &:= \bigcup_{K \in \mathcal{T}_{\mathbf{z}}} K.\end{aligned}$$

For $\mathbf{z} \in \mathcal{N}_{k,0}$, the associated Lagrange basis is given by

$$B_{\mathbf{z}}|_K := \begin{cases} \hat{B}_{\hat{\mathbf{z}}} \circ \phi_K^{-1} & \text{if } \mathbf{z} \in K \in \mathcal{T}_{\mathbf{z}}, \\ 0 & \text{otherwise,} \end{cases}$$

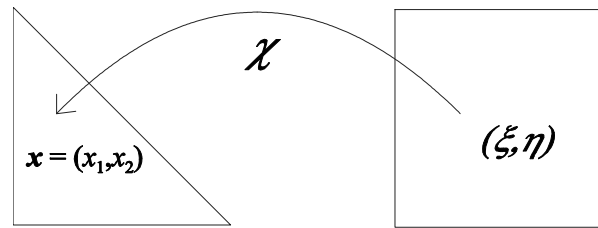
for $\hat{\mathbf{z}} := \phi_K^{-1}(\mathbf{z})$

Computation of the right-hand side

$$r_{\mathbf{z}} = \int_{\Omega} f B_{\mathbf{z}} = \sum_{K \in \mathcal{T}_{\mathbf{z}}} \frac{|K|}{|\hat{K}|} \int_{\hat{K}} \hat{f}_K \hat{B}_{\hat{\mathbf{z}}}$$
$$\hat{f}_K := f \circ \phi_K$$

Employ quadrature formula for the evaluation of the integrals.

Duffy transform



$$\chi(\xi, \eta) := \begin{cases} \xi & d = 1, \\ \xi(1 - \eta_1, \eta_1) & d = 2, \\ \xi(1 - \eta_1, \eta_1(1 - \eta_2), \eta_1\eta_2) & d = 3, \end{cases}$$

$$\det \chi'(\xi, \eta) = \begin{cases} 1 & d = 1, \\ \xi & d = 2, \\ \xi^2 \eta_1 & d = 3. \end{cases}$$

Hence,

$$\int_{\hat{K}} \hat{f}_K \hat{B}_{\hat{\mathbf{z}}} = \begin{cases} \int_0^1 \hat{f}_K \hat{B}_{\hat{\mathbf{z}}} & d = 1, \\ \int_0^1 \int_0^1 \xi \left(\hat{f}_K \hat{B}_{\hat{\mathbf{z}}} \right) \circ \chi \left(\xi, \eta_1 \right) d\eta_1 d\xi & d = 2, \\ \int_0^1 \int_0^1 \int_0^1 \xi^2 \eta_1 \left(\hat{f}_K \hat{B}_{\hat{\mathbf{z}}} \right) \circ \chi \left(\xi, \eta_1, \eta_2 \right) d\eta_2 d\eta_1 d\xi & d = 3 \end{cases}$$

and this can be approximated by Gauss quadrature to high accuracy.

Computation of the system matrix

$$a_{\mathbf{x}, \mathbf{y}} = \int_{\Omega} \langle \nabla B_{\mathbf{z}}, \nabla B_{\mathbf{y}} \rangle = \sum_{K \in \mathcal{T}_{\mathbf{z}} \cap \mathcal{T}_{\mathbf{y}}} \frac{|K|}{|\hat{K}|} \int_{\hat{K}} \langle \mathbf{g}_K \nabla \hat{B}_{\hat{\mathbf{z}}}, \nabla \hat{B}_{\hat{\mathbf{y}}} \rangle$$

with $\mathbf{g}_K := \left(\mathbf{m}_K^{-1} \right)^T \mathbf{m}_K^{-1}$.

Sparsity of the system matrix

Theorem. The system matrix is *sparse*: for any $\mathbf{x}, \mathbf{y} \in \mathcal{N}_{k,0}(\mathcal{T})$ it holds

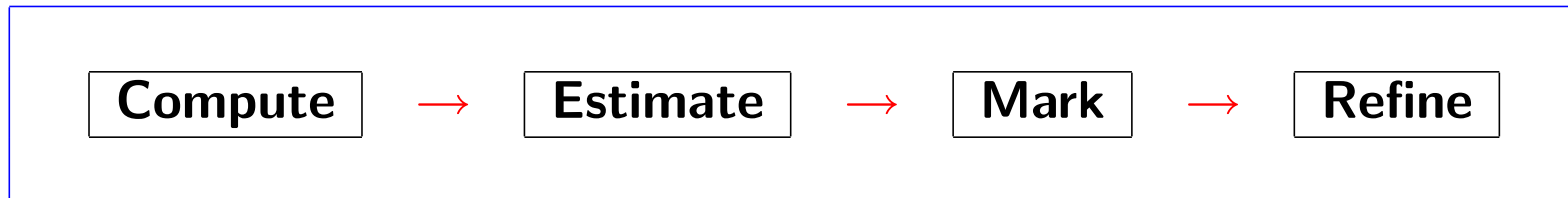
$$\mathcal{T}_{\mathbf{x}} \cap \mathcal{T}_{\mathbf{y}} = \emptyset \implies a_{\mathbf{x},\mathbf{y}} = 0.$$

The number of the non-zero entries per matrix row is bounded by C where C only depends on shape regularity and the polynomial degree p .

Corollary. A matrix-vector multiplication has **linear complexity** $O(N)$ while a general direct solver (Gauss elimination, QR decomposition, ...) of the linear system has a cost of $O(N^3)$. Details \rightarrow *talk of W. Hackbusch*.

Adaptive finite element methods (AFEM)

Idea: Start with a very coarse approximation and set up an algorithm with the structure:



In contrast to *a priori* estimates, *AFEM* uses the computed Galerkin solution to enrich/adapt the space.

Error versus residual:

$$\|u - u_S\|_V = \sup_{v \in V \setminus \{0\}} \frac{a(u - u_S, v)}{\|v\|_V} = \|\mathfrak{R}(u_S)\|_{V'}$$

with *residual* $\mathfrak{R}(u_S) : V \rightarrow V'$ $\mathfrak{R}(u_S)(v) := a(u - u_S, v)$.

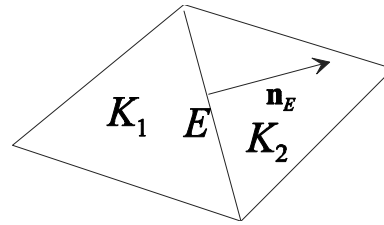
Estimate of residual via local integration by parts (assume $f \in L^2(\Omega)$):

$$\begin{aligned} a(u - u_S, v) &= \int_{\Omega} \langle \nabla(u - u_S), \nabla v \rangle \\ &= - \int_{\Omega} \Delta_{\mathcal{T}}(u - u_S) v + \sum_{K \in \mathcal{T}} \int_{\partial K} \left(\frac{\partial u}{\partial \mathbf{n}_K} - \frac{\partial u_S}{\partial \mathbf{n}_K} \right) v, \end{aligned}$$

where \mathbf{n}_K is the unit outward normal vector for K and $\Delta_{\mathcal{T}}$ denotes the piecewise gradient

$$\Delta_{\mathcal{T}} v|_{\overset{\circ}{K}} = \Delta \left(v|_{\overset{\circ}{K}} \right) \quad \forall K \in \mathcal{T}.$$

Let \mathcal{E}_Ω denote the set of inner element facets (inner mesh points, $d = 1$; inner edges, $d = 2$; inner triangular facets, $d = 3$).



The jump of the normal derivative is given by:

$$\left[\frac{\partial v}{\partial \mathbf{n}_E} \right]_E := \frac{\partial}{\partial \mathbf{n}_E} (u|_{K_2}) - \frac{\partial}{\partial \mathbf{n}_E} (u|_{K_1}).$$

Lemma: If $f \in L^2(\Omega)$ the exact solution u of the Poisson model problem satisfies

$$-\Delta_{\mathcal{T}}u = f \quad \text{and} \quad \left[\frac{\partial u}{\partial \mathbf{n}_E} \right]_E = 0.$$

This leads to

$$a(u - u_S, v) = \int_{\Omega} \text{res}(u_S) v + \sum_{E \in \mathcal{E}_{\Omega}} \int_E \text{Res}(u_S) v.$$

with $\text{res}(u_S) := f + \Delta_{\mathcal{T}}u_S$ $\text{Res}(u_S)|_E := \left[\frac{\partial u_S}{\partial \mathbf{n}_E} \right]_E$.

Remark. The *volume residual* $\text{res}(u_S)$ and the *edge residual* $\text{Res}(u_S)$ are computable.

Employ *Galerkin's orthogonality*: $a(u - u_S, v_S) = 0$ for all $v_S \in S$:

$$a(u - u_S, v - v_S) = \int_{\Omega} \text{res}(u_S)(v - v_S) + \sum_{E \in \mathcal{E}_{\Omega}} \int_E \text{Res}(u_S)(v - v_S).$$

Standard trace estimates with *mesh size function* h lead to

$$a(u - u_S, v) \leq \left(\|h \text{res}(u_S)\|_{L^2(\Omega)} + C_{\text{trace}} \|h^{1/2} \text{Res}(u_S)\|_{L^2(\cup \mathcal{E}_{\Omega})} \right) \times \\ \times \|h^{-1}(v - v_S)\|_{L^2(\Omega)}.$$

Quasi-interpolation

Choose v_S as a “quasi-interpolation” of v with an estimate

$$\|h^{-1}(v - v_S)\|_{L^2(\Omega)} \leq C_{\text{int}} \|v\|_{H^1(\Omega)}$$

Theorem. Let $f \in L^2(\Omega)$ and let u be the exact solution of the Poisson model problem. Let u_S be the Galerkin solution. Then

$$\|u - u_S\|_{H^1(\Omega)} \leq C_{\text{int}} \left(\|h \text{res}(u_S)\|_{L^2(\Omega)} + C_{\text{trace}} \|h^{1/2} \text{Res}(u_S)\|_{L^2(\cup \mathcal{E}_\Omega)} \right).$$

Local error estimator: Define the local error estimator:

$$\eta_K := \sqrt{\|h \operatorname{res}(u_S)\|_{L^2(K)}^2 + \frac{1}{2} \|h^{1/2} \operatorname{Res}(u_S)\|_{L^2(\partial K)}^2}.$$

Then

$$\|u - u_S\|_{H^1(\Omega)} \lesssim \sqrt{\sum_{K \in \mathcal{T}} \eta_K^2}.$$

Marking strategy

Let

$$\eta_{\max} := \max_{K \in \mathcal{T}} \eta_K^2.$$

Choose $\alpha \in]0, 1[$, e.g., $\alpha = 0.7$. Then, mark all simplices for refinement with

$$\eta_K \geq \alpha \eta_{\max}.$$

Remark. AFEM is a strategy to enrich a finite element space by *mesh refinement*.

Outlook

Composite Finite Element spaces are used as a strategy to improve the finite element space *without* increasing the number of unknowns. See next talk.

Thank you for your attention